

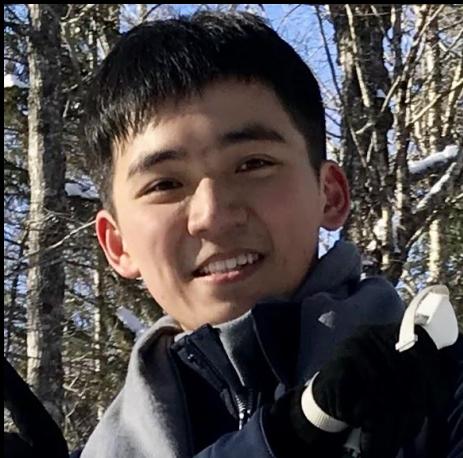
# Hydrodynamics with Helical Symmetry

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# Outline

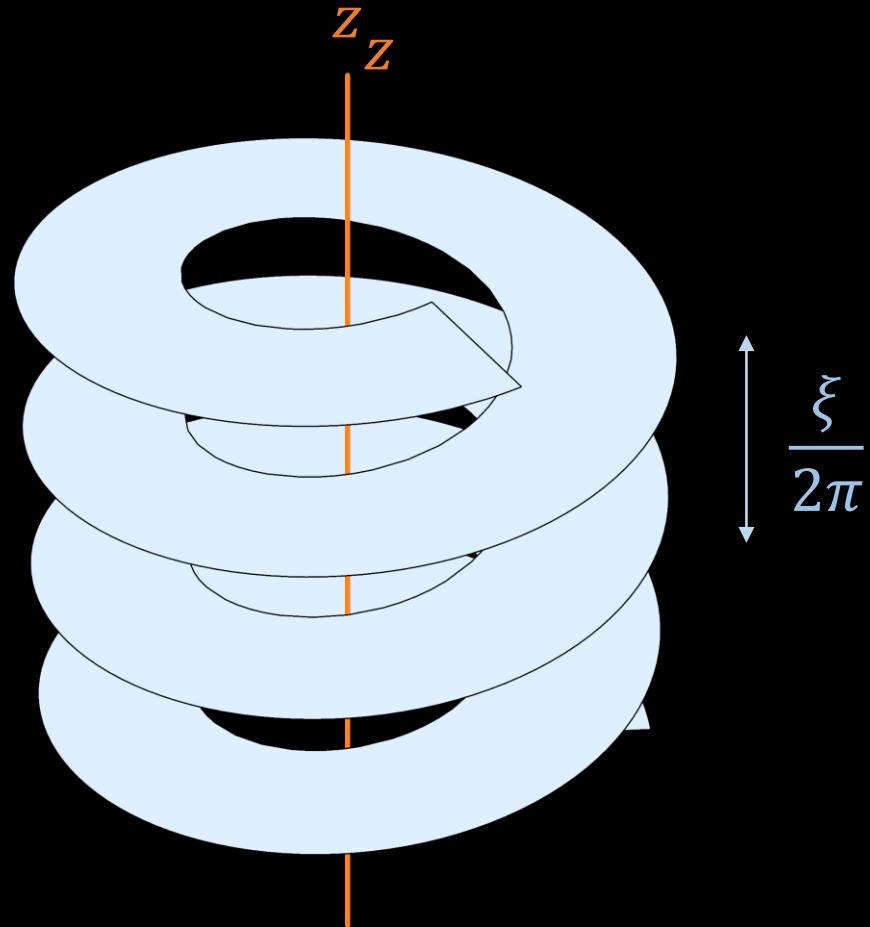
- Helical symmetry
- Review Hydrodynamics
- Helical Hydrodynamics
- Outlook

# Helical Symmetry

- Only a *combination* of translation and rotation is a symmetry
- Corresponding conserved quantity:

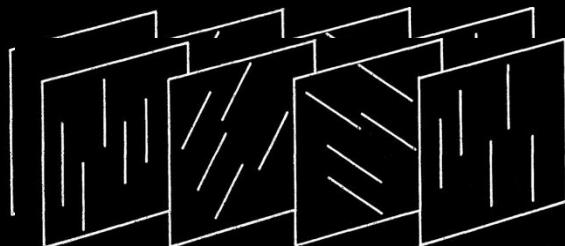
$$K_z = P_z + \frac{L_z}{\xi}$$

- Mixes spacetime symmetries



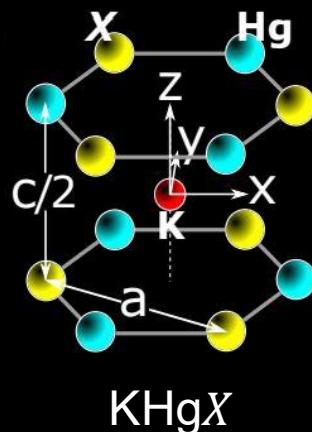
# Realizations:

- Cholesteric liquid crystals (pinned)



[Lubensky, *Hydrodynamics of Cholesteric Liquid Crystals*, 1972]  
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- Non-symmorphic crystal (“discrete helical symmetry”)



[Wang et al, *Hourglass Fermions*, 2016]

# Review of hydrodynamics

- Follow a small number of quantities that are conserved,  $(P_x, P_y, N)$

$$\frac{d}{dt} \int n \, d\mathbf{x} = 0$$

$$\partial_t n + \partial_i J_i = 0,$$

$$\partial_t \pi_A + \partial_i \tau_{iA} = 0.$$

- But what are  $\tau$ ,  $\mathbf{J}$ , and how are they related to  $n, \pi$ ?
- This talk: write down all the terms you can consistent with the symmetries. (Can check within Kinetic Theory and modern EFT framework)

# Gradient expansion

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- Expand  $\tau$  in powers of the hydrodynamic variables and in powers of  $\nabla$  (using  $\mathbf{v}$  instead of  $\boldsymbol{\pi}$ ).
- (linear response is sufficient)

$$\tau_{iA} = A_{iA} n + B_{iAB} v_B + C_{iAk} \partial_k n + D_{iAkB} \partial_k v_B + \dots$$

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- New *dissipationless* terms allowed!

$$C_{iAk} = v_s^2(\alpha \epsilon_{iAz} \delta_{kz} + \beta \epsilon_{iAz})$$

$$\tau_{AB} \sim v_s^2 \alpha \epsilon_{AB} \partial_z n$$

$$\tau_{zA} \sim v_s^2 \beta \epsilon_{AB} \partial_B n$$

# Same game for $\mathbf{J}$ :

$$J_A = n_0 v_A + b \epsilon_{AB} \partial_z v_B$$

$$J_z = n_0 v_z + a \epsilon_{AB} \partial_A v_B$$

- It turns out, for general thermodynamic reasons:

$$a = -\alpha$$

$$b = -\beta$$

# Key Terms in Equations of motion:

$$n_0 \partial_t v_A + v_s^2 \partial_z n + v_s^2 (\alpha + \beta) \epsilon_{AB} \partial_B \partial_z n - \eta \partial_i \partial_i v_A - \eta_z \partial_z^2 v_A = 0$$

$$n_0 \partial_t v_A + v_s^2 \partial_z n + v_s^2 (\alpha + \beta) \epsilon_{AB} \partial_B \partial_z n - \eta \partial_i \partial_i v_A - \eta_z \partial_z^2 v_A = 0$$

$$\partial_t n + n_0 \partial_i v_i - (\alpha + \beta) \epsilon_{AB} \partial_A \partial_z v_B = 0$$

What about  $\pi_z$ ?

Why no *rotational viscosity*  $\eta_o$ ?

What is the value of  $\alpha, \beta$  ?

# What about $\pi_z$ ?

$$\partial_t \pi_z + n_0 \partial_z \mu = -\frac{1}{\xi} (\tau_{xy} - \tau_{yx})$$

- If there is any part of  $\tau_{xy} - \tau_{yx}$  that is *not* a total derivative,  $\pi_z$  is not conserved!

# Rotational viscosity relaxes $\pi_z$

- When angular momentum not conserved:

$$\tau_{xy} - \tau_{yx} = -2\eta_0 \underbrace{(\partial_x v_y - \partial_y v_x)}_{\omega_z}$$

- But a fluid with helix symmetries in equilibrium with density matrix:

$$\rho \sim e^{-\beta(H - \mu N - v_A P_A - w K_z)}$$

- A “boost” in  $K_z$  ( $w \rightarrow w + u$ ) leaves the fluid in equilibrium, so it can not create any entropy!

# Rotational viscosity relaxes $\pi_z$

$$\omega_z \rightarrow \omega_z + 2u/\xi$$

- Not good enough! Instead we need:

$$\tau_{xy} - \tau_{yx} \sim -4\eta_\circ(\omega_z/2 - v_z/\xi)$$

- This gives us the required “not total derivative term”, leading to:

$$\partial_t \pi_z + n_0 \partial_z \mu = -\frac{4\eta_\circ}{\rho \xi^2} \pi_z + \frac{2\omega_z}{\rho \xi}$$

$$\gamma = \frac{4\eta_\circ}{\rho \xi^2}$$

Same relaxation rate as  
calculated in kinetic  
theory and EFT!

$\alpha$  is fixed!

- After  $v_z$  has relaxed:

$$\tau_{xy} - \tau_{yx} \sim -n_0 \xi \partial_z \mu / 2$$

$$\alpha = -n_0 \xi / 2$$

# Summary and Outlook

- Mixing of spacetime symmetries has interesting consequences for hydrodynamics
- Landau, Kinetic Theory, EFT (not presented) all offer useful perspectives on the complications this mixing provides
- “discrete helix”: interesting?
- Other theories that mix multipolar symmetries and translations?

# Thanks a lot for listening!